

What does gravity do with axions?

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arXiv:1307.8024

- the axion as Cold Dark Matter
- the right question: **How to distinguish axions from WIMPs?**
Sikivie's scenario:
 1. “gravitational thermalisation” drives axions to a Bose Einstein condensate
 2. an axion BE condensate supports vortices \Rightarrow caustics in galactic dark matter
- focus on 1.: what is a Bose Einstein condensate?
 - dunno
 - but do need *entropy* generation
 - ? estimate a *dissipative* gravitational interactions for axions in early U?
 - veerry slow...
- anything worth remembering from this talk?

You all know the axion...

- strong CP problem of QCD: $-\frac{1}{4g^2}F_{\mu\nu}^A F^{\mu\nu A} - \frac{\theta}{4g^2}F_{\mu\nu}^A \tilde{F}^{\mu\nu A}$
 θ is CPV — neutron edm $\Rightarrow \theta \lesssim 10^{-10}$
- axial anomaly : rotate θ onto quark masses :(still CPV
- solution: add fields such that rotate θ to the phase of a complex SM-singlet scalar Φ , who gets a vev

$$\Phi \rightarrow f_{PQ} e^{ia/f_{PQ}} \quad f_{PQ} \sim 10^{11} \text{ GeV}$$

Pich deRafael
Pospelov, Ritz

Peccei Quinn

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\Rightarrow only new particle at low-energy is the (pseudo-) goldstone a

$$\text{mixes to pion} : m_a \sim \frac{m_\pi f_\pi}{f_{PQ}} \simeq 6 \times 10^{-5} \frac{10^{11} \text{ GeV}}{f_{PQ}} \text{ eV}$$

$$\text{couplings to SM} \propto \frac{1}{f_{PQ}} \propto m_a$$

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Srednicki NPB85

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light and feebly coupled \Leftrightarrow stellar energy loss bounds: (upper bd on cplg)

$$m_a \lesssim 10^{-2} \text{ eV} \quad (f_{PQ} \gtrsim 10^9 \text{ GeV})$$

Non-thermal axion production cosmology...it is *Cold* Dark Matter!

1. Suppose inflation before Peccei-Quinn Phase Trans.

avoid CMB bounds on isocurvature fluctuations
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* a random in $-\pi \rightarrow \pi$ from one horizon to next,

but $\langle a^2 \rangle_U \text{ today} \sim \pi^2 f_{PQ}^2 / 3$

* one string/horizon

3. QCD Phase Transition ($T \sim 200$ MeV):

* strings go away (radiate cold axion particles, $\vec{p} \sim H \ll m_a$)

Hiramatsu etal 1012.5502

* $m_a(t) : 0 \rightarrow f_\pi m_\pi / f_a$

($\delta\rho \sim \rho$ on horizon-scale !)

when $m_a > H$, axion field oscillates around the minimum...energy density
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Summary:

★ osc. axion field + (cold) axion particles (from strings) redshift like CDM

★ not overclose U for $m_a \gtrsim 10^{-4}$ eV ($f_{PQ} \lesssim 10^{11}$ GeV) recall astrophys: $m_a \lesssim 10^{-2}$ eV
(PQ before inflation can have smaller m_a by tuning $a \ll f_{PQ}$)

★ axion field also grows linear density fluctuations like CDM

Ratra, Hwang+Noh

To cosmologically distinguish axions from WIMPs?

- axions = feebly-coupled, very light one-parameter BSM
 - non-thermal production in cosmology \Rightarrow $\left\{ \begin{array}{l} \text{redshift like WIMPs} \\ \text{grow linear fluctuations like WIMPs} \end{array} \right.$
- \Rightarrow *might axions differ from WIMPs during non-linear structure formation?*
- (Umm... non-linear/N-body is hard!)

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Sikivie's scenario:

Erken, Sikivie, Tam, Yang

1. at $T_\gamma \sim \text{keV}$, “gravitational thermalisation” of axions drives them to a Bose-Einstein Condensate

Saikawa Yamaguchi

2. BEC can support vortices, which allow caustics in the galactic DM distribution
 \Leftrightarrow axion DM signature?

BEC galactic halos:
Rindler-Daller+Shapiro

Rest of the talk is about step 1.

\Rightarrow what is a Bose Einstein condensate?

\Rightarrow what does gravity do with axions during linear structure formation?

\Rightarrow how to find dissipation/thermalisation in gravitational interactions?

What is a Bose Einstein condensate? (I don't know. Please tell me if you do!)

1. in equilibrium stat mech: bosons pile into the $\vec{p} = 0$ mode
2. in equilibrium Finite Temp FT: a phase transition \leftrightarrow form a vev
store a density of conserved charge in a homogeneous + isotropic classical field

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store a density of conserved charge in a homogeneous + isotropic classical field
3. for alkali gases in atomic traps: coherent collective behaviour (all the same \vec{p} ;
but not necc $\vec{p} = 0$)
4. Sikivie says: lowest energy state (not necc homogeneous) pragmatically, it needs to support
vortices?
5. for purposes of this talk: the inhomogeneous classical axion field of misalignment
mechanism is *not* a BE condensate. Need to increase entropy to get there...
 \Rightarrow find dissipation!

*Is a BE condensate just a (non-relativistic) charge-carrying classical field?
Or as well as being “coherent”, does it need to be homogeneous + isotropic, ie,
the $\vec{p} = 0$ mode?*

What does gravity do with axions?

consider early evolution of the Universe, until $\delta\rho \sim \rho$

Eqns of motion inside the horizon thermalisation is causal, so neglecting $H^2/m_a^2, \dots$ for axion field in perturbed FRW (Newtonian gauge, ϕ = Newtonian potential comes from metric) can be obtained from

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$$\ddot{a} + 3H\dot{a} + k^2a + m_a^2a \sim \pm \sum_{q,p} \frac{8\pi G_N m^2}{q^2} a a a + \text{other } \delta\rho \text{ terms}$$

non-linear... can calculate rate for axions to emit a graviton of any wavelength.

Interpretation: the interaction rate from 1 is a thermalisation rate. $\Gamma > H$ after $T_\gamma \sim \text{keV}$, so hugely occupied low- \vec{k} modes equilibrate = form BE condensate.

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2. Or get Eqns of motion for a fluid with scalar perturbations

$$\ddot{\delta} + 2H\dot{\delta} - 4\pi\rho\delta + \frac{c_s^2}{R^2(t)}\nabla^2\delta = 0$$

can solve (in fourier space); gives evolution of axion density fluctuations.

Interpretation: the interaction rate from 1 is a thermalisation rate. $\Gamma > H$ after $T_\gamma \sim \text{keV}$, so hugely occupied low- \vec{k} modes equilibrate = form BE condensate.

But... 2 is the same eqns with different variables, and says (at least some of) those gravitons are growing fluctuations?

Why is that a thermalisation rate? (bath? fluctuations to sum??)

Looking for dissipation in the gravitational interactions of axions

1. Assume BE condensation requires dissipation
2. Assume no dissipation/thermalisation at leading order of classical equations of motion usual non-equilibrium field theory — must sum a bath of fluctuations to dissipate with t -reversal invariant eqns

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4. match axion in perturbed U onto imperfect fluid in FRW:

$$T^i_j(\vec{x}, t) = -\frac{(1 + 2\phi)}{R^2(t)}\partial_i\phi\partial_j\phi = -\eta(t)(\partial_j U^i(\vec{x}, t) + \partial^i U_j(\vec{x}, t))$$

η = viscosity, $U = 4$ -velocity. An imperfect fluid can grow density fluctuations, but contains dissipation...

5. estimate a dissipation scale:
< the Jeans length $1/\sqrt{m_a H}$, distance below which fluctuations oscillate due to axion pressure

Summary (review of well-known things)

Maybe New Physics is *light!* Maybe the axion, feebly coupled solution to the θ -problem, is the Dark Matter. It should satisfy 10^{-4} eV $\lesssim m_a \lesssim 10^{-2}$ eV (upper bound from astrophys).

Non-thermally produced axions can be the *Cold* Dark Matter of the Universe. There are axion field oscillations, and if PQPT after inflation, cold axion particles from strings.

To confirm/distinguish axions from WIMPs?

1– direct detection: find WIMPs or axions in terrestrial searches (CAST, ADMX...)

2– during structure formation: redshift, grow linear density growth same

\Rightarrow *are axions different from WIMPs during non-linear structure formation ?*

Speculations...on distinguishing axions from WIMPs during non-linear structure formation

? Maybe a classical field forms a galaxy differently from WIMPs?

⇒ ask a friendly N-body person to write an axion code?

⇒ (if there are axions from strings, do they condense to a field??)

Sikivie answer :

1-gravity drives Bose Einstein condensation of axions

2-the BE condensate supports vortices, and the vortices give observable caustics

caveat about 1: remains to be shown that the gravitational interaction rate of axions is a thermalisation rate, or how it changes the axion distribution

* leading order classical equations (no entropy generation?) for axions in perturbed FRW reproduce the gravitational interaction rate Sikivie identifies as a thermalisation rate. Maybe its the gravitational interactions growing the density fluctuations?

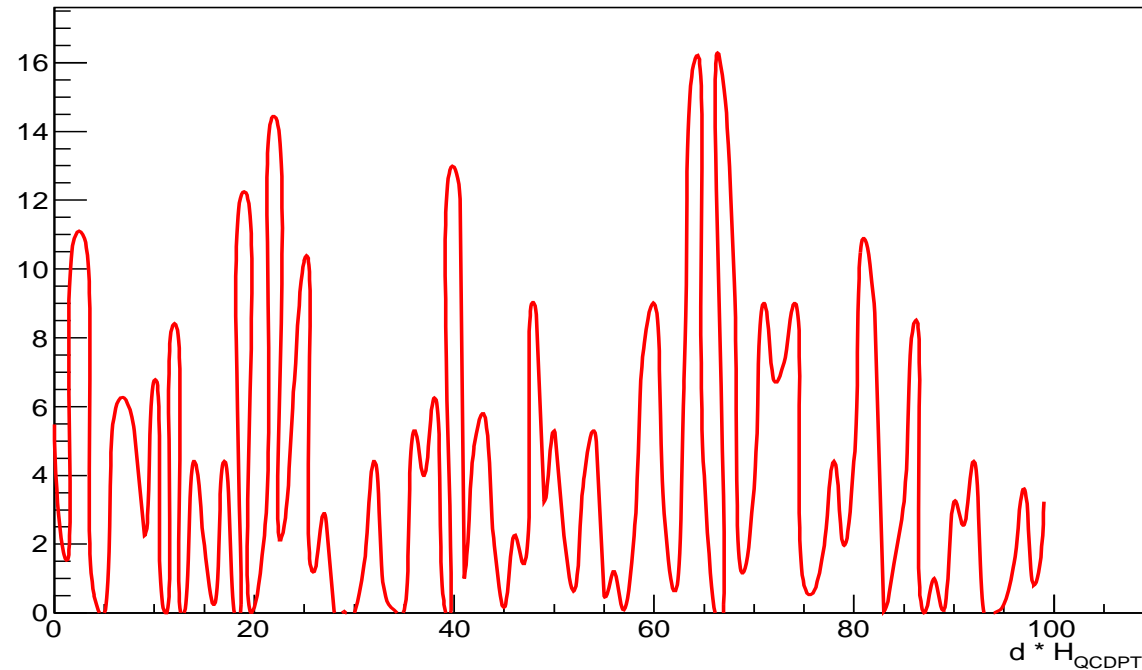
* We found some dissipative gravitational interactions, but suppressed by p_a^2/m_a^2 ...

Backup

Inhomogeneities are $\mathcal{O}(1)$ on the QCD horizon scale

$a(\vec{x}, t)$ random from one horizon ($\sim 5\text{km}$) to next; $\rho_a(\vec{x}, t) \simeq m_a^2 a^2(\vec{x}, t)$

axion density at the QCDPT

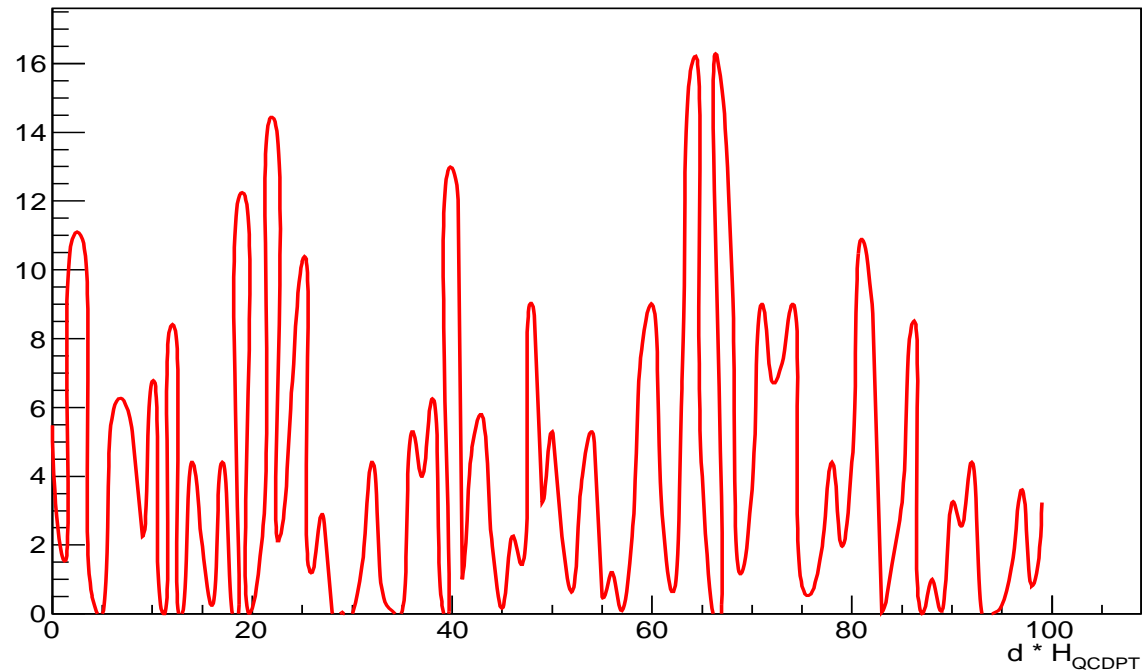


\Rightarrow its *not* a spatially homogeneous distribution of particles various momenta

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axion density at the QCDPT



But how can axions form a *homogeneous-on-QCD-horizon-scale* bose-einstein condensate = zero mode of field? ??

$v = H_{QCDPT}/m_a \lesssim 10^{-6}c$...not “free-stream” QCD-horizon distance before t_{eq} :

$$d(t) = \int^t \frac{H_{QCDPT}}{m_a R(t')} dt' \sim \frac{H_{QCDPT}}{m_a} \frac{1}{H(t)R(t)} = \frac{R(t)}{m_a} \ll \frac{R(t)}{H_{QCDPT}}$$

(RD U, $R(t) = 1 @ QCDPT$)

The (beautiful) calculation of Saikawa and Yamaguchi

Suppose PQ PT after inflation. The classical axion field can be represented as a coherent state of axion particles (of momentum $\lesssim H_{QCDPT}$).

QFT rate for axions (momentum \vec{k}) to emit gravitons:

$$i\frac{\partial \hat{n}_k}{\partial t} = [\hat{H}_{int}, \hat{n}_k] \simeq \frac{G_E}{H(t)^2} \rho_a^2 \gg H(t) n_k$$

Saikawa+Yamaguchi

(evaluated in coherent state \Leftrightarrow classical field caln.)

Sikivie interprets as gravitational thermalisation rate: hugely occupied low- \vec{p} modes, equilibrium after $T_\gamma \lesssim \text{keV}$, \rightarrow BE condensate.

But are some of those gravitons expanding the U, and some growing fluctuations?

Why is that a thermalisation rate??

thermalisation in closed unitary systems?

$$\text{entropy} = \sum_{\text{states } s} P_s \ln P_s \quad \text{increases}$$

- unitary evolution creates no entropy \Leftrightarrow *NO* entropy generation in closed systems
... *BUT*... can calculate “effective” thermalisation: a subset of observables evolve towards equilibrium expectations
 \Rightarrow the “rest” of the system is the bath??
- ex: couple two SHOs. Solve one, substitute into Eqns of second, and find dissipation.
- ... $K - \bar{K}$ evolution is non-unitary, because not also follow 2π 3π states...

? \Rightarrow divide axions+gravity into

1. U expansion + structure growth
2. other fluctuations which are the bath?

gravity and the second law

1. undergraduate memories say that gravitational collapse of a gas cloud to a star respects the second law...
2. story of $\Omega_{baryon} = 1$ U
 - (a) quasi-homogeneous dust clouds collapse
 - (b) ...generations of stars, supernovae, black holes...
 - (c) proton decays...
 - (d) venerable homogeneous and isotropic U full of photons and gravitons
3. so gravitational thermalisation of axions will happen.
But does it happen before the U a year old?

Particles vs fields

Develop field operator

$$\hat{a}(t, \vec{x}) = \frac{1}{[R(t)L]^{3/2}} \int \frac{d^3k}{(2\pi)^3} \left\{ \hat{b}_{\vec{k}} \frac{\chi(t)}{\sqrt{2\omega}} e^{i\vec{k}\cdot\vec{x}} + \hat{b}_{\vec{k}}^\dagger \frac{\chi^*(t)}{\sqrt{2\omega}} e^{-i\vec{k}\cdot\vec{x}} \right\}$$

then write the coherent state:

$$|a(\vec{x}, t)\rangle \propto \exp \left\{ \int \frac{d^3p}{(2\pi)^3} a(\vec{p}, t) b_{\vec{p}}^\dagger \right\} |0\rangle$$

which satisfies $\hat{b}_{\vec{q}} |a(\vec{x}, t)\rangle = a(\vec{q}, t) |a(\vec{x}, t)\rangle$ (can check $\hat{b}_{\vec{q}} \{1 + \int \frac{d^3p}{(2\pi)^3} a(\vec{p}, t) b_{\vec{p}}^\dagger\} |0\rangle = a(\vec{q}, t) |0\rangle$)

where the classical field is

$$a(t, \vec{x}) = \frac{1}{[R(t)L]^{3/2}} \int \frac{d^3k}{(2\pi)^3} \left\{ a(\vec{k}, t) \frac{\chi(t)}{\sqrt{2\omega}} e^{i\vec{k}\cdot\vec{x}} + a^*(\vec{q}, t) \frac{\chi^*(t)}{\sqrt{2\omega}} e^{-i\vec{k}\cdot\vec{x}} \right\}$$