Global Higgs Fits

and implications for extended Higgs sectors

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Based on work with G. Belanger, B. Dumont, U. Ellwanger, J.F. Gunion, et al.



outline



- Introduction & motivation
- Using the experimental results
 - Disentangling production modes
 - Combined signal strengths
- Global coupling fits
 - constraints on reduced couplings
 - constraints on invisible/unseen decays
 - interplay with dark matter searches
- Implications for specific models
 - 2HDM, inert Higgs, triplets, SUSY
- Conclusions

"On the presentation of the LHC Higgs results," arXiv:1307.5865

"Global fit to Higgs signal strengths and couplings and implications for extended Higgs sectors," **arXiv:1306.2941**

"Status of invisible Higgs decays," arXiv:1302.5694

"LHC constraint on light neutralino dark matter in the MSSM," **arXiv:1308.3735**

and work in progress

motivation

- The 2012 discovery of a SM-like Higgs at a mass around 125 GeV is a first triumph for the LHC physics program.
- The data collected at $\sqrt{s} = 7$ and 8 TeV already provide quite a comprehensive picture of the production and decay properties of the 125 GeV Higgs boson
- However, while the Higgs completes our picture of the SM, it still leaves many fundamental questions open (naturalness, hierarchy problem)
 → new physics beyond the SM ?
- Given the absence of direct signals for BSM at the LHC, the Higgs data and their interpretation currently provide the crucial guidelines for BSM theories.



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The Higgs as guide to BSM ?

- In BSM theories, the Higgs production cross sections, decay branching ratios, kinematic distributions, and even the number of Higgs particles may differ from SM predictions.
- Have to distinguish between two classes of models by whether or not the selection efficiencies and detector acceptances for the various channels are independent of the model parameters.
- In the former case, SM-like tensor structure of the 125 GeV Higgs
 → signal strength modifiers
- Otherwise perform MC simulation, apply cuts, compare with exp. data
 → ideally want fiducial cross sections
- See suggestions "On the presentation of the LHC Higgs results" in **1307.5865**.



using the experimental results

signal strengths

- Presenting results in terms of $\mu = \sigma/\sigma_{SM}$ is a very convenient way to quantify agreement with SM expectations.
- BSM contributions affect production as well as decay rates → detailed breakdown in terms of production×decay modes needed to test (B)SM.
- Experimental practice: data related to a single decay mode H→Y are divided into categories (or "sub-channels") *I*, in order to improve sensitivity or discrimination among the production mechanisms.
- Need to interpret this in terms of "pure" production modes: ggF, VBF, VH, ttH
 - → need efficiencies!





signal strengths in "theory space"



The likelihood in terms of $\mu(X,Y)$ allows for reinterpretation of the results in models where the efficiency and acceptance for each (X,Y) is approximately unchanged with respect to the SM.

using sub-channel information

• The likelihood in terms of $\mu(X,Y)$ can be approximately recomputed combining the χ^2 of all categories I using an efficiency-weighted sum:

$$\mu_{I}(Y) = \sum_{X} \mu(X, Y) T(I, X) \sigma(X_{\text{SM}}) \text{BR}(H_{\text{SM}} \to Y)$$
selection efficiencies for each production mode, normalized to one.

Jack calls this the "transfer matrix".

- It is critical that for each of the categories *I* the selection efficiencies (and uncertainties thereon) be provided for all production modes. Unfortunately this is not yet done in a systematic way :-(
- NB important correlations may be missed in this approach, e.g., from migration of events between categories.

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| Expected signal and estimated background | | | | | | | | | _ | | |
|--|----------------------------|-------|--------|-------------------------|---------|----------|-----------------------|-----------|--------------|------------------|---|
| Event classes | | | SM Hig | ggs boso | on expe | cted sig | nal ($m_{\rm H}$ =12 | T | | | |
| Event classes | | | | 2.521 100-000-000-00 | 10 | 23 | $\sigma_{ m eff}$ | FWHM/2.35 | МЕТ | | CMS preliminary |
| | 0.00000 | Total | ggH | VBF | VH | ttH | (GeV) | (GeV) | Electron | _ | √s = 7 TeV, L = 5.1 fb ⁻¹ (MVA) √s = 8 TeV, L = 19.6 fb ⁻¹ (MVA) |
| TeV 5.1 fb ⁻¹ | Untagged 0 | 3.2 | 61.4% | 16.8% | 18.7% | 3.1% | 1.21 | 1.14 | Muon | > - | Event Class |
| | Untagged 1 | 16.3 | 87.6% | 6.2% | 5.6% | 0.5% | 1.26 | 1.08 | Di-jet loose | ר <mark>ש</mark> | Combined m _H = 125.0 GeV |
| | Untagged 2 | 21.5 | 91.3% | 4.4% | 3.9% | 0.3% | 1.59 | 1.32 | Di-jet tight | | σ/σ _{SM} = 0.78+0.28-0.26 |
| | Untagged 3 | 32.8 | 91.3% | 4.4% | 4.1% | 0.2% | 2.47 | 2.07 | Untagged 3 | ₩ - | |
| 1 | Dijet tag | 2.9 | 26.8% | 72.5% | 0.6% | — | 1.73 | 1.37 | Untagged 2 | | |
| 8 TeV 19.6 fb ⁻¹ | Untagged 0 | 17.0 | 72.9% | 11.6% | 12.9% | 2.6% | 1.36 | 1.27 | Untagged 1 | - | - |
| | Untagged 1 | 37.8 | 83.5% | 8.4% | 7.1% | 1.0% | 1.50 | 1.39 | Untagged 0 | | + - |
| | Untagged 2 | 150.2 | 91.6% | 4.5% | 3.6% | 0.4% | 1.77 | 1.54 | Di-jet | > | _ |
| | Untagged 3 | 159.9 | 92.5% | 3.9% | 3.3% | 0.3% | 2.61 | 2.14 | Untagged 3 | 0 - | |
| | Dijet tight | 9.2 | 20.7% | 78.9% | 0.3% | 0.1% | 1.79 | 1.50 | Untagged 2 | ⊢ → | _ |
| | Dijet loose | 11.5 | 47.0% | 50.9% | 1.7% | 0.5% | 1.87 | 1.60 | Untagged 1 | | |
| | Muon tag | 1.4 | 0.0% | 0.2% | 79.0% | 20.8% | 1.85 | 1.52 | Untagged 0 | | |
| | Electron tag | 0.9 | 1.1% | 0.4% | 78.7% | 19.8% | 1.88 | 1.54 | -10 | -5 0 | 5 10 |
| | $E_{\rm T}^{\rm miss}$ tag | 1.7 | 22.0% | 2.6% | 63.7% | 11.7% | 1.79 | 1.64 | | | Best Fit σ/σ _{SM} |

T(I,X) a good example

from CMS-PAS-HIG-13-001 ($H \rightarrow \gamma \gamma$, mass-fit MVA analysis)

but unfortunately not yet available for all channels from both ATLAS and CMS (also not for CMS $H \rightarrow \gamma \gamma$ CiC analysis)

using sub-channel information

 Reconstruction of 68 and 95% CL contours from sub-channel info (black/grey) and comparison to official ATLAS/CMS results (blue)



arXiv:1307.5865

• NB important correlations may be missed in this approach. In particular, some systematic uncertainties lead to migration of events between categories, and these uncertainties can dominate over the statistical ones.

$\mu(ggF+ttH)$ vs $\mu(VBF+VH)$ plots

 It has become standard that for each decay mode the experiments present 68% and 95% CL contours in the μ(ggF+ttH) versus μ(VBF+VH) plane:

CMS Preliminary $\sqrt{s} = 7 \text{ TeV}, L \le 5.1 \text{ fb}^{-1} \sqrt{s} = 8 \text{ TeV}, L \le 19.6 \text{ fb}^{-1}$ µ _{VBF,VH} $\mu_{VBF+VH} \times B/B_{SM}$ 10 $H\to\tau\tau$ ATLAS Preliminary $H \rightarrow WW$ $vs = 7 \text{ TeV}: \int Ldt = 4.6-4.8 \text{ fb}^{-1}$ $H \rightarrow ZZ$ $\sqrt{s} = 8 \text{ TeV}: \int Ldt = 13-20.7 \text{ fb}^{-1}$ + H → bb + $H \rightarrow \gamma \gamma$ Standard Mode $ZZ^{(1)} \rightarrow 4I$ $\rightarrow WW^{(*)} \rightarrow hh = -$ 2 2 0 0 -2 m_H = 125.5 GeV 2 6 8 0 2 3 4 5 1 -1 3 $\mu_{ggH,ttH}$ $\mu_{ggF+ttH} \times B/B_{SM}$

- This is a boon for interpretation studies because the fundamental production modes are already "unfolded" from the experimental categories.
- Could be extended to other $\mu(X,Y)$ vs $\mu(X',Y')$ combinations, e.g. WH, ZH for $H \rightarrow bb$

VH=WH+ZH

$\mu(ggF+ttH)$ vs. $\mu(VBF+VH)$: limitations

- μ (VBF+VH) assumes custodial symmetry.
- If only 68% and 95% CL contours are given, one first needs to reconstruct the likelihood. Simplest solution is fitting a 2D Gaussian:



Gaussian fits (red/orange contours) to signal strengths in the $\mu(ggF+ttH)$ vs $\mu(VBF+VH)$ plane and comparison to official ATLAS and CMS results (in blue).

> In each case, we approximately reconstruct the likelihood by fitting a bivariate normal distribution to the 68% CL contour given by the collaboration

> > Scalars 2013, Warsaw

It would be of great advantage to have the full likelihood information in the μ(ggF+ttH) vs μ(VBF+VH) plane ... or other relevant planes



 Preferably this information should be directly available in numerical form (via INSPIRE → DOI → searchable and citable)

a big step forward



This week the ATLAS collaboration has taken an important step forward by making the likelihood function for three key measurements about the Higgs available to the world digitally. [K. Cranmer, Quantum Diaries, 12-Sep-2013]



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| Information Files | NSPIRE |
|--|---|
| Data from Figure 7 from: Measurements of Higgs production and couplings in diboson final states v ATLAS detector at the LHC | boson with the |
| ATLAS Collaboration (Aad, Georges (Freiburg U.) []) Show all 2923 authors | <u>iors</u> |
| Cite as: ATLAS Collaboration (2013) HepData, http://doi.org/10.7484/INSPIREHEP.DATA.A78C.HK44 | |
| Description: -2 log Likelihood for the H $\rightarrow \gamma\gamma$ channel in the (µ_ggF+ttH * B/BSM, B/BSM) plane for a Higgs boson mass mH = 125.5 GeV. Preview not available | , µ_VBF+VH * Figure 7 |
| Note: * Temporary entry * This dataset complements the following publication: <u>Measurements of Higgs boson production and couplings in diboson final states</u> <u>ATLAS detector at the LHC</u> Record created 2013-09-11, last modified 2013-09-11 | $M_{H} = 125.5 \text{ GeV}$ |
| | $\mu_{ggF+ttH} \times B/B_{SM}$ |

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combined signal strengths

Combining ATLAS, CMS and Tevatron results

Fitting 2D Gaussians to the 68% CL contours from the experiments, we construct a combined likelihood in the (ggF+ttH,VBF+VH) plane for each final state:



arXiv:1306.2941

Combined µ's



w/o Tevatron

| | $\hat{\mu}(ggF + ttH)$ | $\hat{\mu}(\text{VBF} + \text{VH})$ | ρ |
|----------------------------|------------------------|-------------------------------------|--------|
| $\gamma\gamma$ | 0.98 ± 0.28 | 1.72 ± 0.59 | -0.38 |
| VV | 0.91 ± 0.16 | 1.01 ± 0.49 | -0.30 |
| $b\overline{b}/	au	au$ | 0.98 ± 0.63 | 0.97 ± 0.32 | -0.25 |
| au	au | 1.07 ± 0.71 | 0.94 ± 0.65 | -0.47 |
| $\overline{b}\overline{b}$ | -0.23 ± 2.86 | 0.97 ± 0.38 | 0 |



Agrees frustratingly well with SM :-(but, there's still room for sizable deviations ...

global fits

arXiv:1306.2941

Coupling Fits

• Need to specify the Lagrangian

$$\mathcal{L} = g \left[C_V \left(M_W W_\mu W^\mu + \frac{M_Z}{\cos \theta_W} Z_\mu Z^\mu \right) - C_U \frac{m_t}{2M_W} \bar{t}t - C_D \frac{m_b}{2M_W} \bar{b}b - C_D \frac{m_\tau}{2M_W} \bar{\tau}\tau \right] H.$$

C's scale couplings relative to SM ones; $C_U=C_D=C_V=1$ is SM.

- Couplings to gluons and photons: we compute $\overline{C_g}$ and $\overline{C_Y}$ from C_U , C_D , C_V ; we also allow additional loop contributions ΔC_g and ΔC_Y from new particles $\rightarrow C_g = \overline{C_g} + \Delta C_g$ and $C_Y = \overline{C_Y} + \Delta C_Y$
- Calculation of σ×BR following the recommendations of the LHC Higgs Cross Section Working Group, arXiv:1209.0040
- Fit includes ATLAS, CMS and Tevatron results from Moriond and LHCP 2013.
- NB when relevant we also include searches for invisible decays. In particular ATLAS ZH→II+MET gives B(inv)<0.65 at 95% CL.



$\Delta C_{g}, \Delta C_{Y}$ Fit



68%, 95% and 99.7% CL contours

Best fit: $\Delta C_g = -0.06$ $\Delta C_Y = 0.13$

$$\chi^2_{\rm min}/{\rm dof} = 0.84 \ (SM: 0.82)$$

$C_U, C_D, C_V \text{ Fit} \qquad (\Delta C_g = \Delta C_Y = 0)$



Comparison with ATLAS and CMS coupling fits



invisible decays



invisible decays



invisible decays



constrained by searches for invisible H

unseen decays



 In principle all the Higgs production*decay rates can be kept fixed by scaling up the C's while adding a new, unseen decay mode with branching ratio B_{new}.

• For
$$C = C_U = C_D = C_V : C^2 = 1/(1 - B_{new})$$

D. Zeppenfeld et al, hep-ph/0002036A. Djouadi et al, hep-ph/0002258M. Duhrssen et al, hep-ph/0406323

• This gives a flat direction in C_U , C_D , C_V . For $C_V \le I$ however, we can still get a strong constraint on B_{new} similar to the case of invisible decays. At 95% CL:

i) $\mathcal{B}_{\text{new}} < 0.21$ for a SM Higgs with allowance for unseen decays; *ii)* $\mathcal{B}_{\text{new}} < 0.31$ for C_U, C_D free, $C_V \leq 1$ and $\Delta C_{\gamma} = \Delta C_g = 0$; and *iii)* $\mathcal{B}_{\text{new}} < 0.39$ for $C_U = C_D = C_V = 1$ but $\Delta C_{\gamma}, \Delta C_g \neq 0$ allowed for.

total width



dotted: C_U, C_D, C_V \leq I dashed: C_U, C_D, C_V solid: C_U, C_D, C_V, Δ C_g, Δ C_Y

testing custodial symmetry

Fit to ATLAS and CMS results as in arXiv:1306.2941 but taking C_W and C_Z as independent parameters. $C_{WZ} = C_W / C_Z$



Best fit: C_Z=1.1, C_W=0.98

[internship J. Bernon]

extended Higgs sectors

inert doublet model

SM plus a second Higgs doublet \tilde{H}_2 which is odd under a Z_2 symmetry \rightarrow DM $V = \mu_1^2 |H_1|^2 + \mu_2^2 |\tilde{H}_2|^2 + \lambda_1 |H_1|^4 + \lambda_2 |\tilde{H}_2|^4 + \lambda_3 |H_1|^2 |\tilde{H}_2|^2 + \lambda_4 |H_1^{\dagger} \tilde{H}_2|^2 + \frac{\lambda_5}{2} \left[\left(H_1^{\dagger} \tilde{H}_2 \right)^2 + \text{h.c.} \right]$



$$\lambda_{L,S} = \frac{1}{2}(\lambda_3 + \lambda_4 \pm \lambda_5)$$

$$m_h^2 = \mu_1^2 + 3\lambda_1 v^2,$$

$$m_{\tilde{H},(\tilde{A})}^2 = \mu_2^2 + \lambda_{L(S)} v^2,$$

$$m_{\tilde{H}^{\pm}}^2 = \mu_2^2 + 0.5\lambda_3 v^2.$$

$$\frac{\lambda_3}{2} = \frac{1}{v^2} \left(m_{\tilde{H}^+}^2 - m_{\tilde{H}}^2 \right) + \lambda_L$$

inert doublet model



Scalars 2013, Warsaw

mssm with light neutralinos

 Scan over weak-scale MSSM parameter space, requiring consistence at 95% CL with flavor constraints, Higgs mass and h signal strengths in all channels, as well as DM relic density and DM direct searches
 SK, Kulkarni, Laa, Lessa, Proschofsky-

Spindler, Waltenberger, in preparation]

• SUSY mass limits from LHC via "SModelS" Simplified Model approach



conclusions

- mH~125 GeV is quite fortunate: we can detect the Higgs in many different channels.
- ATLAS and CMS measurements already provide quite a comprehensive picture
 - \rightarrow test for deviations from SM predictions
 - \rightarrow constrain models of new physics
- Publication of µ likelihoods by ATLAS is an important step towards maximizing impact and utility of LHC results

"I truly hope that this becomes standard practice for the LHC." (Kyle)

• Eventually we need to go beyond 2D μ 's :

$\mathcal{L}(m_H, \mu_{\text{ggF}}, \mu_{\text{ttH}}, \mu_{\text{VBF}}, \mu_{\text{ZH}}, \mu_{\text{WH}})$

 For testing models with e.g. new tensor structure, we need to go beyond signal strengths → fiducial cross sections



[>]robability density

500 1000 1500 2000 2500 3000 $\widetilde{\chi}_{\star}^{\pm}$ mass [GeV]

an exciting way ahead (to unmapped territory, I hope)



backup

On the presentation of the LHC Higgs results

F. Boudjema, G. Cacciapaglia, K. Cranmer, G. Dissertori, A. Deandrea, G. Drieu la Rochelle, B. Dumont, U. Ellwanger, A. Falkowski, J. Galloway, R.M. Godbole, J.F. Gunion, A. Korytov, S. Kraml, H.B. Prosper, V. Sanz, S. Sekmen

Abstract:

We put forth conclusions and suggestions regarding the presentation of the LHC Higgs results that may help to maximize their impact and their utility to the whole High Energy Physics community.

Conclusions and suggestions from the workshops "Likelihoods for the LHC Searches", 21-23 Jan 2013 at CERN, "Implications of the 125 GeV Higgs Boson", 18-22 March 2013 at LPSC Grenoble, and from the 2013 Les Houches "Physics at TeV Colliders" workshop.

signal strengths beyond 2D

- Eventually, we want to test ggF, ttH,VBF, ZH and WH separately, which means that we need a more detailed break down of the channels beyond 2D plots.
- The optimum would of course be to have the full statistical model available
 → RooFit workspaces ?
- What we would like to advocate (as a compromise) is that for each final state Y the experiments give the signal strength likelihood in the 6D form

$$\mathcal{L}(m_H, \mu_{
m ggF}, \mu_{
m ttH}, \mu_{
m VBF}, \mu_{
m ZH}, \mu_{
m WH})$$

- This way, a significant step could be taken towards a more precise fit in the context of a given BSM theory.
- The likelihood could be communicated either as a standalone computer library or as a large grid data file.
- Open point: final state correlations \rightarrow covariance matrix ?

additional Higgses

- In searches for additional Higgs states φ, the contributions from the SM-like Higgs boson at ~125 GeV should be treated like any other SM background.
- Results should always be reported as bounds on $\sigma \times BR$ for any ϕ !



• Degenerate states at ~125 GeV are a special issue \rightarrow Andrey

S. Kraml

summary of recommendations

For each Higgs decay mode $Y(\gamma\gamma, WW, ZZ, b\bar{b}, \tau\tau$ are currently considered) provide the likelihood \mathcal{L} of the signal strengths in the $(\mu(X, Y), \mu(X', Y))$ plane, as shown in Figure 3. The grouping X = ggF + ttH, X' = VBF + VH is well motivated, but additional choices of X and X' should be considered when appropriate for the given analysis. The content of the plots should *always* be provided also in numerical form, *e.g.*, as a ROOT file or as a simple text file with a grid. In addition to the combined results, results should also be given separately for each \sqrt{s} .

To go a step further and overcome the limitations induced by 2D projections and/or combining production modes, provide the signal strength likelihood as a function of m_H , separated into all five production modes ggF, ttH, VBF, ZH and WH; *i.e.* for each decay mode considered give the likelihood in the 6D form $\mathcal{L}(m_H, \mu_{ggF}, \mu_{ttH}, \mu_{VBF}, \mu_{ZH}, \mu_{WH})$. Ideally, this should again also be done separately for each \sqrt{s} .

summary of recommendations

Whenever possible, provide kinematic event selection criteria that can approximately be reproduced by phenomenologists, *e.g.*, using Monte Carlo event generators.² The desired information is: the complete cut flow, estimated number of background events, expected event yields for all the SM Higgs processes, and the observed number of signal events or limits thereon. For MVA-based analyses, it would be of great value if a simplified version of the final MVA could be given.

Concerning searches for additional Higgs-like states with masses above or below 125 GeV, provide the results including the injection of a signal with the properties of a SM-like Higgs boson at 125–126 GeV. Moreover, always present the results as bounds on pure ($\sigma \times B$) in addition to any model interpretation.

summary of recommendations

In addition to direct model-dependent interpretations of data, the long-term goal should be to develop a consistent scheme for publishing fiducial cross sections ($\sigma^{\text{fid}} \times \mathcal{B}$), either measurements or limits for null search results, as done conventionally for SM processes.

We suggest that this supplementary material is made available via INSPIRE [50]. This way the complete set of information will be searchable, citable, and accessible from a single point.

NB the document is open to discussion.

 \rightarrow your input can help to extend and improve it.