

# Naturalness in General Gauge Mediation

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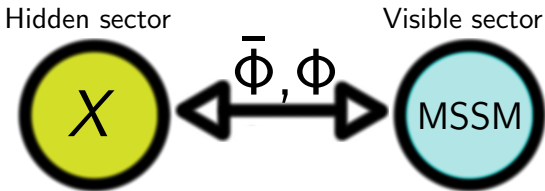
based on: Z. Lalak and ML, arXiv:1302.6546

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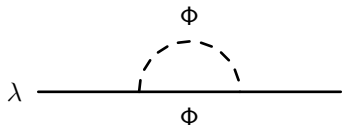


# SUSY breaking mediation

- Supergravity
  - No control over mixing between families  $\rightarrow$  large FCNC
- Gauge mediation
  - SUSY is spontaneously broken  $\rightarrow$  singlet  $\langle X \rangle = X + \theta^2 F$
  - breaking is transmitted through messengers  $W = \lambda \bar{\Phi} X \Phi$
  - messengers  $\bar{\Phi}, \Phi$  interact with MSSM fields only via gauge interactions

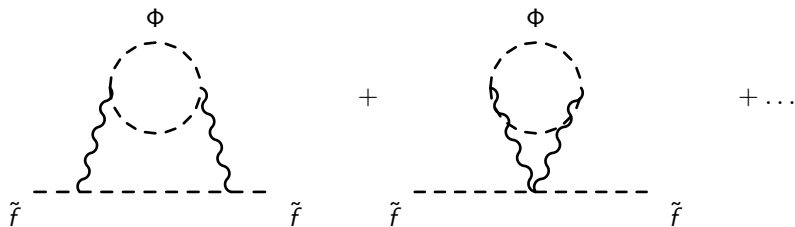


# Gauge mediated soft terms



A Feynman diagram showing a fermion line (solid line) with a loop of scalar particles (dashed line) labeled  $\Phi$ . The loop is connected to the fermion line at two points. The diagram is followed by an arrow pointing to the equation  $M_i = \frac{\alpha_i F}{4\pi X}$ .

$$\lambda \text{ --- } \text{loop} \text{ --- } \lambda \implies M_i = \frac{\alpha_i F}{4\pi X}$$



Two Feynman diagrams showing fermion self-energy corrections. The first diagram shows a fermion line (dashed line) with a loop of scalar particles (dashed line) labeled  $\Phi$ . The second diagram shows a fermion line (dashed line) with a loop of scalar particles (dashed line) labeled  $\Phi$ . The diagrams are separated by a plus sign, and the second diagram is followed by a plus sign and an ellipsis.

$$\tilde{f} \text{ --- } \text{loop} \text{ --- } \tilde{f} \quad + \quad \text{loop} \text{ --- } \tilde{f} \quad + \dots$$

$$\implies m_{\tilde{f}}^2 = 2 \sum_i C_i(f) \left( \frac{\alpha_i}{4\pi} \right)^2 \left| \frac{F}{X} \right|^2$$

Meade, Shih and Seiberg 0801.3278

Gauge mediated soft terms can be expressed by just six parameters

- Three gaugino masses

$$M_1 = \frac{\alpha_1}{4\pi} m_Y, \quad M_2 = \frac{\alpha_2}{4\pi} m_w, \quad M_3 = \frac{\alpha_3}{4\pi} m_c,$$

- Three parameters determining scalar masses  $\Lambda_c^2$ ,  $\Lambda_w^2$ ,  $\Lambda_Y^2$  which give

$$m_{\tilde{f}}^2 = 2 \left[ C_3(f) \left( \frac{\alpha_3}{4\pi} \right)^2 \Lambda_c^2 + C_2(f) \left( \frac{\alpha_2}{4\pi} \right)^2 \Lambda_w^2 + C_1(f) \left( \frac{\alpha_1}{4\pi} \right)^2 \Lambda_Y^2 \right],$$

- Only negligible A-terms are generated.

Two specific models

Carpenter et al. 0805.2944

- GGM1

$$W_{GGM1} = X_i(y^i \bar{Q}Q + r^i \bar{U}U + s^i \bar{E}E),$$

with three independent parameters  $\Lambda_Q, \Lambda_U, \Lambda_E$

- GGM2

$$W_{GGM2} = X_i(y^i \bar{Q}Q + r^i \bar{U}U + s^i \bar{E}E + \lambda_q^i \tilde{q}q + \lambda_l^i \tilde{l}l),$$

with five independent parameters  $\Lambda_Q, \Lambda_U, \Lambda_E, \Lambda_q, \Lambda_l$

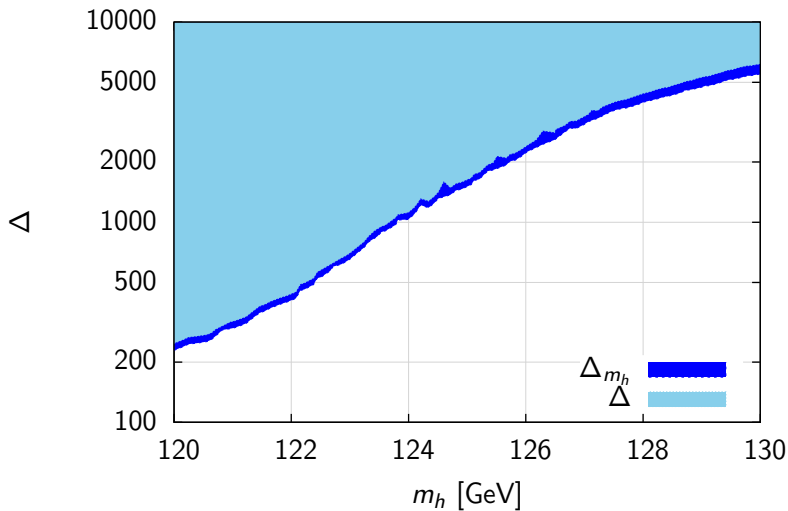
- fine-tuning from parameter  $a$

$$\Delta_a = \left| \frac{\partial \ln m_Z^2}{\partial \ln a} \right|.$$

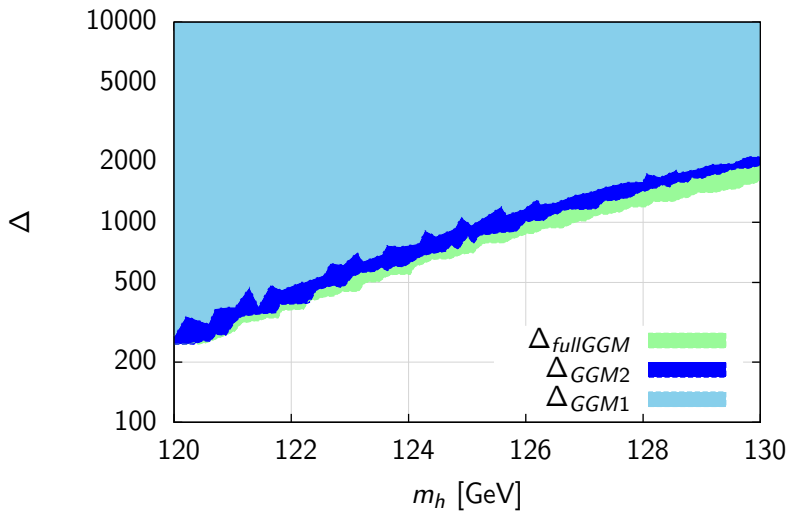
- fine-tuning coming from a whole set of parameters  $a_i$

$$\Delta = \max_{a_i} \Delta_{a_i}.$$

# FT in mSUGRA



# FT in GGM





## reducing fine-tuning

Assuming that parameters are not independent of each other, but instead are functions of some fundamental parameters. For example, if gaugino masses  $M_i$  are given functions of parameter  $M_{\frac{1}{2}}$  we obtain

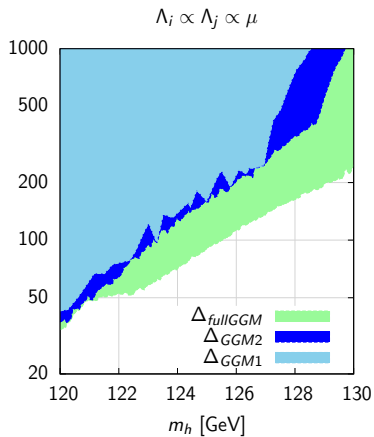
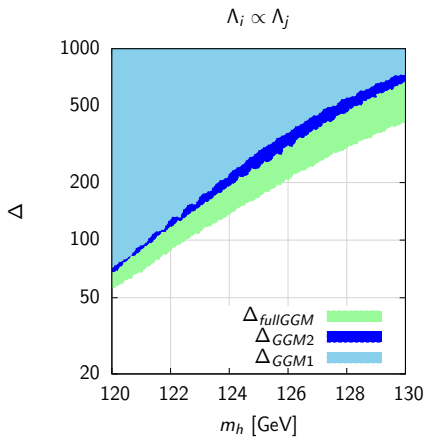
$$M_i = f_i(M_{\frac{1}{2}}),$$
$$\Delta_{M_{\frac{1}{2}}} = \left| \frac{\partial \ln M_Z^2}{\partial \ln M_{\frac{1}{2}}} \right| = \left| M_{\frac{1}{2}} \frac{f'_i(M_{\frac{1}{2}})}{f_i(M_{\frac{1}{2}})} \frac{\partial \ln M_Z^2}{\partial \ln M_i} \right|.$$

If  $f_i$  are simply proportional to  $M_{\frac{1}{2}}$  one finds

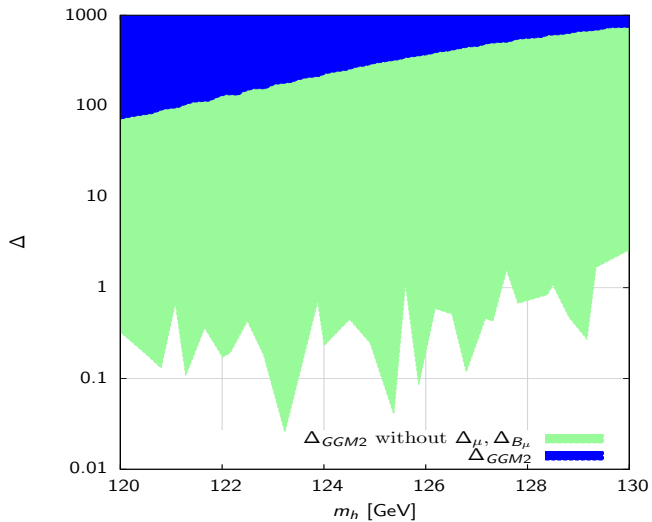
$$\Delta_{M_{\frac{1}{2}}} = \left| \sum_{i=1}^3 \frac{\partial \ln M_Z^2}{\partial \ln M_i} \right|.$$

If these functions were logarithms

$$M_i(M_{\frac{1}{2}}) = \tilde{m} \ln \frac{M_{\frac{1}{2}}}{Q}, \quad \Delta_{M_{\frac{1}{2}}} = \left| \sum_{i=1}^3 \frac{\tilde{m}}{M_i} \frac{\partial \ln M_Z^2}{\partial \ln M_i} \right|.$$



# fine-tuning from only gauge mediated soft terms



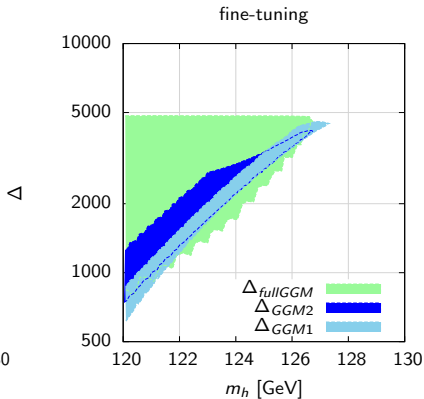
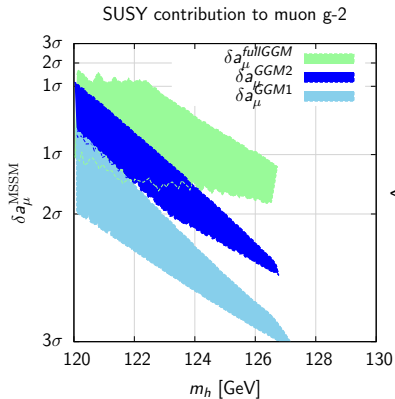
- discrepancy between measurement and SM prediction:

$$\delta a_\mu = a_\mu^{\text{EXP}} - a_\mu^{\text{SM}} = (2.8 \pm 0.8)10^{-9}.$$

- The simplest approximation of SUSY contribution

$$\delta a_\mu^{\text{SUSY}} \approx \left( \frac{g_1^2 - g_2^2}{192\pi^2} + \frac{g_2^2}{32\pi^2} \right) \frac{m_\mu^2}{M_{\text{SUSY}}^2} \text{tg } \beta,$$

Problem: We need heavy superpartners ( $M_{\text{SUSY}}$ )



- ① GGM predicts smaller fine-tuning than mSUGRA
- ② for  $m_h = 126\text{GeV}$  fine-tuning always larger than 100 unless one includes only gauge mediated soft terms
- ③ including  $g_\mu - 2$  raises fine-tuning about four times, but its still possible to obtain  $g_\mu - 2$  within  $1\sigma$  bound
- ④ decrease of the Higgs mass down to 123 GeV reduces the fine-tuning by a factor of 2.